The iron oxidation state of garnet by electron microprobe: Its determination with the flank method combined with major-element analysis

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ABSTRACT

We have established a method to determine the oxidation state of Fe using the electron microprobe—the “flank method.” This method has the possibility to analyze simultaneously the major-element chemistry on the same spot. The method is based on Fe L X-ray emission spectra that show a concomitant change of both the intensity and the wavelength of the soft Fe Lα and Fe Lβ emission lines. The method has been developed for garnets using well-characterized synthetic samples, and an empirical correction scheme has been established for self-absorption of the Fe L X-ray emission. This enables the unambiguous correlation of measured intensity ratios of Fe Lα/Fe Lβ X-ray emission to the oxidation state in garnet. The flank method can be used with routine microprobe settings (e.g., 15 kV, 60 nA) and, for garnets, with minimum beam size. This way, Fe3+/ΣFe and major and trace elements can be determined within 10 minutes on the microscale. The resulting accuracy of Fe3+/ΣFe is then ±0.04 for garnets at 10 wt% total Fe. Application to natural garnet megacrysts with 8 wt% total Fe and Fe3+ known from Mössbauer spectroscopy yields excellent agreement between the two methods. For practical application, the calibration can be achieved by using three natural standards and fitting the data to Fe3+ and Fe2+ with linear equations. Through repeated measurements on homogeneous samples, the error in Fe3+/ΣFe becomes approximately ±0.02 (1σ). Optimization of analytical parameters such as beam current, measurement time, and repetition of analyses, will further lower the uncertainty.

Keywords: Iron oxidation state, flank method, self-absorption, electron microprobe, garnet

INTRODUCTION

The difference in the geochemical behavior of Fe2+ and Fe3+ in conjunction with the Fe2+/Fe3+ ratio is an important indicator of the oxygen fugacity during mineral formation. It makes the determination of the Fe oxidation state in minerals a highly sought after aim in earth sciences. Specifically, the in situ speciation of Fe on the microscale is relevant to geology and material sciences, since geosystems and synthetic materials generally are of small grain size or zoned. Techniques for the determination of Fe oxidation state with micrometer or nanoscale spatial resolution have gained new impetus due to recent advances in X-ray spectroscopy with either electron beam sources [electron microprobe (EPMA) or transmission electron microscopy (TEM)] or with X-ray beam sources (synchrotron facilities).

The best spatial resolution for Fe3+/ΣFe determination is currently achieved with TEM using energy-loss spectroscopy (EELS) (Garvie and Craven 1994; Golla and Putnis 2001; van Aken and Liebscher 2002; van Aken et al. 1998, 1999). EELS opens the possibility of quantitative nanometer-scaled valence state mapping (Golla and Putnis 2001). For quantitative measurements, a spot size of about 1 μm in the TEM diffraction mode is used to avoid electron beam-induced oxidation (van Aken and Liebscher 2002). Absolute errors for Fe3+/ΣFe with EELS then range between ±0.02 and ±0.04 depending upon which evaluation method is used with regard to energy-loss in near-edge structures. X-ray photoelectron spectroscopy [XPS, e.g., Raeburn et al. (1997a, 1997b)] and X-ray absorption near edge structure (XANES) spectroscopy (Bajt et al. 1994; Berry et al. 2003; Cressey et al. 1993; Delaney et al. 1998; Dyr et al. 2002; O’Neill et al. 2006; Schmid et al. 2003; Wilke et al. 2001) take an intermediate position in terms of the two-dimensional spatial resolution for Fe3+/ΣFe determination. XANES now allows Fe speciation with a 10 × 15 μm X-ray beam size and a spectrum acquisition time of 20–30 min with an error in Fe3+/ΣFe between ±0.03 and ±0.10 for garnets (Dyr et al. 2002). A very recent technique is RIXS (resonant inelastic X-ray scattering), which enables a detailed analysis of the Fe K pre-edge spectral features that are sensitive to the oxidation state in conventional absorption spectroscopy (de Groot 2001; Glatzel and Bergmann 2005).

The technique most widely used at present for Fe3+/ΣFe determination in earth sciences is Mössbauer spectroscopy. Although it is a bulk method it has gained importance for micrometer-sized areas through the development of the “milliprobe” (McCammmon et al. 1991, 2001). Nonetheless, a spatial resolution of >50 μm is large compared to the resolution of an electron microprobe and long acquisition times are restrictive for routine use.

Electron-microprobe techniques with a spatial resolution on the order of 1 μm are based either on the change of the intensity or on the change of the wavelength of the soft Fe Lα and Fe Lβ X-ray emission lines. Both have been exploited in earlier attempts, generally only with semi-quantitative results [for an overview
see Höfer et al. (1994, 2000)]. This semi-quantitative nature likely stems from several factors: (1) neither the wavelength shift nor the change in \(L_\beta/L_\alpha\) intensity ratios taken alone are very sensitive; (2) self-absorption of the Fe\(_{\alpha}\) X-ray emission line by Fe was not accounted for, i.e., a correction for total Fe was missing; and (3) different mineral species were not treated separately. Attempts to correct the \(L_\beta/L_\alpha\) data for self-absorption with matrix-correction programs have also failed (Armstrong 1999; O’Nions and Smith 1971) since mass-absorption coefficients are highly uncertain in the proximity of absorption edges, particularly near the \(L\) edges (Duncumb 1993). More recently, Fialin et al. (2004, 2001) analyzed the Fe\(_{\alpha}\) peak shift of several structurally different minerals or glasses with either Fe\(_{\alpha}^2\) or Fe\(_{\alpha}^3\) alone and took the overall difference between their positions as a normalizing measure. This difference increases with an increasing Fe content. In principle, this could yield a quantitative method as self-absorption is accounted for except for the fact that: (1) uncertainties in the peak position measurements are about 15% at 10 wt% Fe [compare Fig. 3 in Fialin et al. (2001)], and (2) data on mixed-valency minerals are lacking, which necessitates the unproven assumption that there is a linear increase of Fe\(_{\alpha}^2/\Sigma\)Fe between both curves for pure Fe\(_{\alpha}^2\)- and pure Fe\(_{\alpha}^3\)-bearing minerals.

Our approach in evaluating the Fe\(_{\alpha}\) and Fe\(_{\beta}\) X-ray emission lines is the “flank method” (Höfer et al. 1994). It is a hybrid method that consists of both measuring the \(L_\beta/L_\alpha\) intensity ratios and the peak shift. This results in a dramatic increase in sensitivity. The flank method has been applied successfully to sodic amphiboles by Enders et al. (2000) and to wüstites (Höfer et al. 2000). First results on synthetic garnets were presented by Höfer et al. (1994) and by Höfer (2002) in a paper focused on the correlation between flank method data, Mössbauer quadrupole splitting, and crystallographic data.

In the present work, it will be shown that the flank method yields precise and accurate determinations of Fe\(_{\alpha}^3/\Sigma\)Fe in garnets and that it can be combined with simultaneous quantitative elemental determination. These advances become possible by extending our studies to a large number of synthetic, compositionally different garnets. Their analysis enabled us to establish systematic variations to correct for self-absorption as a dominating factor for Fe\(_{\alpha}\) X-ray emission lines, as already pointed out by Höfer et al. (1994). Since we have carried out flank-method measurements using both Jeol and Cameca microprobes, we can provide potential users with a practical overview concerning the steps involved in flank-method measurements for either electron microprobe (see Appendix).

### SAMPLE SELECTION

All samples used were synthetic garnets with simple cation exchange vectors on either the octahedral or dodecahedral site (Table 1). They were synthesized at various temperature and pressure conditions in hydrothermal, belt, and multi-anvil apparatus and were donated by the persons whose names are listed in Table 1. The cation-substitution scheme on the dodecahedral or octahedral sites in these garnets is also given, together with the range in Fe\(_{\alpha}^3/\Sigma\)Fe. Samples with Fe\(_{\alpha}^2\), Mg\(_{\alpha}^2\), Mn\(_{\alpha}^2\), Ca\(_{\alpha}^2\) on the dodecahedral site and Fe\(_{\alpha}^2\), Al\(_{\alpha}^3\), Cr\(_{\alpha}^3\) on the octahedral site cover the full range from 0 to 1. The Fe\(_{\alpha}^3/\Sigma\)Fe ratio in pyrope-

#### ELECTRON PROBE MICROANALYSIS (EPMA)

Measurements have been carried out over the years first with a Camebax microprobe at the University of Maine, then with a Cameca SX 50 at TU München and TU Darmstadt, and now with the Jeol JXA 8900RL at the University of Frankfurt. Our Jeol microprobe is equipped with a turbo molecular pump to minimize oil vapor in the column. The other microprobes have oil diffusion pumps. Earlier results on a part of the sample set have been already published (Höfer 2002). These measurements were re-evaluated in the course of this study with our new method for self-absorption correction, and the samples were also newly measured with the combination of flank method and major element analysis.

The flank-method measurements carried out with the Cameca microprobes were obtained by accumulating 2 7 700 counts at one position each of the Fe\(_{\beta}\) and L\(\beta\) peak flanks at 706.4 and 716.3 eV, respectively. This corresponds to an error of \(\pm 0.6\)% due to counting statistics as described by Höfer et al. (1994, 2000). Measurement conditions were 10 kV, 60 nA, TAP crystal, and scanning mode (approx. 5 \(\times\) 3 \(\mu\)m) or defocused point mode (3 \(\mu\)m). For separate quantitative elemental analysis, we used 15 kV, 40 nA, and point mode. Oxygen was also analyzed with the almandine standard of McGregor et al. (1992).

For the combined flank method and quantitative elemental analysis using the Jeol microprobe, the Fe\(_{\beta}\) and L\(\beta\) intensities were measured for a fixed time of up to 600 s each, typically 300 s. This has the disadvantage of different counting statistics for \(L_\alpha\) and for L\(\beta\), but the advantage of a fixed measurement time of 10 min per analyzed spot, even at very low Fe concentrations. One spectrometer with a TAP crystal and smallest detector slit (300 \(\mu\)m) was used for flank-method measurements. With the remaining 4 spectrometers, up to 12 elements (Si, Ti, Al, Cr, Fe, Mg, Ca, Mn, Ni, Na, K, P) were measured simultaneously, whereby counting times (in s) on peak and background were: Ni 60/60, Ca 60/30 (CaSiO\(_3\)); Ti 60/60 (MnTiO\(_3\)); Al 60/60 (Al\(_2\)O\(_3\)); Cr 60/30 (Cr\(_2\)O\(_3\)); Fe 60/30 (Fe\(_3\)SiO\(_4\)); Mg 60/60 (MgSiO\(_3\)); Mn 60/60 (MnO); Ni 60/30 (NiO); Na 20/20 (albite); K 30/15 (K\(_2\)O\(_3\)PO\(_4\)). The accelerating voltage was 15 kV with a 60 nA beam current. The beam in point mode was either defocused to 3 \(\mu\)m or focused to the minimum beam diameter achievable with the W filament.

#### THE FLANK METHOD

### Derivation of flank measuring positions from difference and absorption spectra

In Figure 1a, the Fe\(_{\beta}\) X-ray emission lines for two garnets [almandine and andradite] containing Fe\(_{\alpha}^2\) and Fe\(_{\alpha}^3\), respectively, are plotted vs. energy. Peak maxima are shifted to higher energies for andradite (Fe\(_{\alpha}^3\)) compared to almandine (Fe\(_{\alpha}^2\)), and the

#### TABLE 1

<table>
<thead>
<tr>
<th>System</th>
<th>Cation exchange</th>
<th>(\text{Fe}_{\alpha}^3/\Sigma)Fe</th>
<th>Reference</th>
</tr>
</thead>
<tbody>
<tr>
<td>Almandine-Pyrope</td>
<td>Mg(<em>{\alpha}^2) \leftrightarrow Fe(</em>{\alpha}^3)</td>
<td>0</td>
<td>Geiger and Rossman 1994</td>
</tr>
<tr>
<td>Almandine-Pyrope</td>
<td>Mg(<em>{\alpha}^2) \leftrightarrow Fe(</em>{\alpha}^3) + Ca(_{\alpha}^2)</td>
<td>0</td>
<td>Geier et al. 1987</td>
</tr>
<tr>
<td>Almandine-Pyrope</td>
<td>Mg(<em>{\alpha}^2) \leftrightarrow Fe(</em>{\alpha}^3)</td>
<td>0</td>
<td>G.P. Brey (unpubl)</td>
</tr>
<tr>
<td>Almandine-Spessartine</td>
<td>Fe(<em>{\alpha}^2) \leftrightarrow Mn(</em>{\alpha}^2)</td>
<td>0</td>
<td>Geier and Feenstra 1997</td>
</tr>
<tr>
<td>Almandine-Skiagite</td>
<td>Al(<em>{\alpha}^3) \leftrightarrow Fe(</em>{\alpha}^3)</td>
<td>0 – 0.4</td>
<td>Woodland and O’Neill 1993</td>
</tr>
<tr>
<td>Andradite-Skiagite</td>
<td>Ca(<em>{\alpha}^2) \leftrightarrow Fe(</em>{\alpha}^3)</td>
<td>0.4 – 1</td>
<td>Woodland and Ross 1994</td>
</tr>
<tr>
<td>Andradite-Grossular</td>
<td>Fe(<em>{\alpha}^2) \leftrightarrow Al(</em>{\alpha}^3)</td>
<td>1</td>
<td>A.B. Woodland (unpubl)</td>
</tr>
</tbody>
</table>
intensity ratios of $L\alpha$ to $L\beta$ are also higher. In principle, there are three way to exploit this behavior: (1) measure the peak shift; (2) evaluate the change in peak intensities either by measuring the intensities at the peak maximum or by measuring the peak areas; or (3) determine the ratio $L\beta$/L$\alpha$ on the flanks of $L\beta$ and $L\alpha$ at the defined energy positions of the minima and maxima of the difference spectrum (Fig. 1a). This distinctly enhances the sensitivity compared to the first two methods because it is a hybrid fortifying both effects ("flank method," Höfer et al. 1994 and 2000).

The measurement positions for the flank method were determined from absorption spectra that may be calculated from the X-ray emission spectra at different accelerating voltages as logarithmic intensity ratios (Liefeld 1968; Fischer and Baun 1968). Almandine and andradite X-ray emission spectra at 2.5 and 10 kV were collected and used to derive ln($I_{3\alpha}/I_{\alpha}$) (Figs. 1b and 1c). For better comparison, the spectra were normalized to equal intensities at the flank-method measurement position for Fe$\beta$. It is quite apparent and most pronounced in almandine that the $L\alpha$ X-ray emission maxima are shifted to lower energies at 10 kV due to absorption on the high-energy side of the peak. The calculated absorption spectra are shown in the lower parts of Figures 1b and 1c. The $L_3$ absorption edge overlaps exactly with the $L\alpha$ flank-method measuring position and moves across this measurement position as the Fe$^{3+}$ content of the sample increases. Thus, by determining the intensities at this position, one directly monitors the effect of the valence-dependent movement of the absorption edge. This is the most probable reason for the high sensitivity of the flank method.

By comparing Figures 1a and 1b, it can be seen that the absorption spectrum calculated for Fe$^{3+}$ (almandine) closely resembles the difference spectrum between andradite and almandine. This demonstrates that the differences between the andradite and almandine X-ray emission spectra are mainly caused by self-absorption. It also shows that the calculation of a difference spectrum is an easy and effective tool to determine the flank-method measurement positions.

**Advantage of the flank method over “conventional” $L\beta$/L$\alpha$ ratios**

From the literature, $L\beta$/L$\alpha$ intensity ratios measured at the peak maximum are available for natural garnets (Albee and Chodos 1970; Pavičević et al. 1989; Rösler 1989). These data are plotted in Figure 2 and show a clear dependency of this ratio on Fe$^{3+}$/ΣFe, but the scatter is quite large. We analyzed several members of the almandine-"skiagite" and "skiagite"-andradite solid solution series in a similar way but used peak area ratios. The results with respect to resolution are fairly similar (Fig. 2). The application of the flank method yielded a substantial increase of the range of ratios from andradite to "skiagite" to almandine, but the relationship is not linear. A few pyrope-almandine garnets with only Fe$^{2+}$ were also measured and it is clear that both our data and those from the literature (Pavičević et al. 1989) fall off the correlation lines.

**Reproducibility, contamination, error sources**

The reproducibility of the flank-method measurements was determined using a synthetic almandine sample (GV2, Geiger and Rossman 1994) and the $L\beta$/L$\alpha$ ratio was monitored in 72 measurements over 14 months (Fig. 3). We obtained a mean value of 1.454 ± 0.018 (1σ). This value is prone to errors by faulty conditions such as spectrometer mispositioning, contamination, or major changes in laboratory conditions. Figure 3 shows one drastic deviation of the $L\beta$/L$\alpha$ ratio beyond the usual standard deviation as the result of a room temperature decrease of about 2°C due to an overnight thunderstorm with a 15°C temperature drop. Almandine is therefore used as a monitor for daily checks of the microprobe settings.

Contamination by stray hydrocarbons in the column may be a problem for the long wavelengths of the FeL lines. Tests were carried out by repeated measurements on the same spot of an almandine sample at 10 kV, 60 nA over a scanned area of 3
\* Flank Method 15 kV
\* Flank Method 10 kV
\* Peak Method 10 kV
\* Literature data 15 kV

\* Peak Area Method 10 kV

\* Literature data 15 kV (Peak Maximum)

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**FIGURE 2.** \(L_\beta/L_\alpha\) intensity ratios of natural garnets measured at 15 kV at peak maximum vs. \(Fe^{3+}/\Sigma Fe\) (data from literature, see text). Also shown are our 10 and 15 kV data on synthetic garnets as peak-area and as flank-method ratios. The peak-area ratios correlate in a way similar to the shift in peak maximum (literature) with only a small resolution. The resolution of the flank-method ratios is distinctly higher.

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**FIGURE 3.** Reproducibility of the flank-method ratio of almandine monitored over several months. A single outlier occurred as a result of a 2°C room temperature drop in the microprobe lab. The 1σ deviation of the mean value is indicated by the shaded area.

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**FIGURE 4.** (a) Contamination tests with repeated measurements on the same spot of an almandine sample using an air jet (10 kV, 60 nA, electron beam on a scanned area of \(3 \times 5 \mu m\), Cameca SX 50). Without an air jet, the intensity of \(FeL_\beta\) decreases significantly after 30 min, while \(FeL_\alpha\) remains constant. (b) The intensity ratio \(L_\beta/L_\alpha\) decreases accordingly after 30 min. The intensity ratio remains within the 99% confidence level for \(\approx 1.5\) h with the use of an air jet.

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**RESULTS AND THE INTRODUCTION OF THE SELF-ABSORPTION CORRECTION**

Systematic variation of \(L_\beta/L_\alpha\) flank method ratios in synthetic garnets—results at 10 kV

Figure 6 shows all flank-method results obtained at 10 kV plotted against the total Fe content. The data (mean values generally consist of 10 measurements) span a quadrilateral where \(Fe^{3+} = 0\) (I), \(Fe^{3+} = constant\) and equal to 3 in the structural formula.

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\(\times 5 \mu m\) in a Cameca SX50 microprobe. After about 30 min of measurement time, the \(FeL_\beta\) intensity gradually decreased while \(FeL_\alpha\) is not affected (Fig. 4a). Accordingly, \(L_\beta/L_\alpha\) also gradually decreases (Fig. 4b). An explanation could be the beginning of oxidation in the sample that would reduce foremost the intensity of \(L_\beta\) (Tossell et al. 1974). Using an air jet as an anticontamination device, the \(FeL_\beta\) intensity decrease is less pronounced. This adversity points to the effect of contamination. After 1.5 h, the \(L_\beta/L_\alpha\) ratio remained within the 99% confidence level (Fig. 4). Therefore, the use of an air jet or other anticontamination device does not seem necessary if measurement times of more than 30 min on one spot are avoided. Nevertheless, it is advisable to measure \(L_\beta\) first and \(L_\alpha\) second since contamination does not seem to affect the \(L_\alpha\) intensity at all.

For further tests on reproducibility, stability in the electron beam, contamination, and different measurement conditions for low Fe garnets, we used a garnet megacryst from the Monastery kimberlite in South Africa (Mon B) with 9.0 wt% total Fe. Different spot sizes from “0” \(\mu m\) (= minimum with the W filament, approximately 0.2 \(\mu m\) at 60 nA) to defocused (1 or 2 \(\mu m\)) were combined with 30 or 300 s measurement times for \(L_\beta\) (measured first) and \(L_\alpha\), with repetitions on the same sample spot or with “grid” measurements on new spot positions. The measurements were taken over a period of 4 weeks. As can be seen from Figure 5, all these different combinations yielded the same mean value for the \(L_\beta/L_\alpha\) intensity ratio. One small increase in the 1 \(\mu m\) beam diameter is due to a room temperature increase of 0.3°C during these measurements. These results show that there is no dependence on the beam size or on the number of repetitions, that garnets are very stable in the electron beam, and that there is no contamination problem. The scatter for the short 30 s measurement times is by nature larger than for the 300 s measurements. However, the data show a Gaussian distribution around the average value both for the “grid” measurements with 399 different points and for the 100 repeated measurements on the same spot.
FIGURE 5. Reproducibility tests with the icel microprobe at 15 kV, 60 nA on a garnet megacryst (Mon B) with 9 wt% Fe. The data are grouped according to combinations of different measuring times (30 s and 300 s both on Lβ and Lα), different beam sizes ("0 μm" = minimum, 1 and 2 μm), repetitions on the same sample spot or "30 s grid" measurements on different positions for each measurement. The gray bar indicates the ±1σ uncertainty of the mean Lβ/Lα value for 10 measurements (600 s each) on 10 different spots (data not shown). This corresponds to the normal measuring routine.

(II), Fe³⁺ = constant and equal to 2 in the structural formula (III), and Fe²⁺ = 0 (IV). Pyrrole-almandine samples that seemingly were not consistent with the other data in Figure 2 now form line I in Figure 6. Also plotting along this line are all other samples with only Fe²⁺, but with varying amounts of Mg, Mn, Ca, and with Cr replacing Al. In the almandine-"skiagite" solid-solution series, Fe³⁺ replaces Al³⁺ on the octahedral site. In pure "skiagite," the octahedral site is filled with Fe³⁺ and the dodecahedral site with Fe²⁺. The Fe³⁺/ΣFe ratio is 0.4. This ratio increases from "skiagite" to andradite and reaches one in andradite. Although Fe³⁺ remains constant per formula unit, the total Fe content is changed due to the replacement of Fe²⁺ by Ca²⁺. In the andradite-grossular garnets, Fe³⁺ exchanges with Al³⁺. Examination of this diagram shows that Lβ/Lα strongly depends on the Fe²⁺ content (lines I and III). Thus, self-absorption of the generated Fe L X-ray emission by the Fe atoms in the sample must cause the variation in Lβ/Lα. The slope of line I in Figure 6 directly yields the self-absorption effect of Fe²⁺ on the measured ratio since total Fe equals Fe²⁺ in these garnets. The behavior of group III garnets shows the same dependence on the total Fe content, but it is shifted parallel because of the higher total Fe content that in turn is a consequence of Fe³⁺ in these samples.

In Figure 7, the data for the group I–IV garnets are plotted against their Fe²⁺ content alone (as known from Mössbauer spectroscopy or determined from their stoichiometry). These data were collected at 10 kV, 60 nA with the Cameca SX50 in Darmstadt without simultaneous quantitative analysis. A linear relationship arises with a correlation coefficient of R = 0.99 (Fig. 7), onto which all garnets fall regardless of their chemical compositions or their Fe³⁺/Fe²⁺ to ferric Fe ratio. It should be noted, however, that the scatter in Lβ/Lα is also due to the inhomogeneity of some of these synthetic samples. This problem can be overcome by simultaneously analyzing the chemical composition in the same spot (see below). At 0 wt% Fe²⁺, the Lβ/Lα value from the fit is 0.501, which agrees perfectly with the theoretical value of 0.5 predicted for the electronic transition probabilities from the M⁴,5 orbital to the L₂ and L₃ energy levels of the Fe atom as a result of the occupancy of these orbitals. This is shown later in Figure 18. This value remains unchanged by varying amounts of Fe³⁺. We conclude that Lβ/Lα is mainly governed by a linear self-absorption effect caused by Fe²⁺ that is superimposed by a second self-absorption effect due to the total Fe with a small contribution stemming from Fe³⁺. This can be used to quantify Fe³⁺/ΣFe, as will be demonstrated below.
Effect of self absorption: Enhancing the flank method sensitivity

Self-absorption is the dominant factor that affects $L\beta/L\alpha$. This enables us to quantify Fe$^{3+}$ contents using the flank method. Self-absorption depends on the Fe concentration but also on the path length of the Fe L X-rays within the sample. An increased path length results from a greater penetration depth of the electrons that in turn depends on the accelerating voltage. The consequence is an increased interaction volume as shown schematically in Figure 8. The increase in the accelerating voltage from 10 to 15 and 20 kV increases $L\beta/L\alpha$ and the difference of the andradite and almandine ratios, so that sensitivity is enhanced (Fig. 9). At 2.5 kV, on the other hand, the differences are minimal since self-absorption is negligible.

Plotting the data vs. the total Fe content (Fig. 10) shows no dependency of the ratio at 2.5 kV. Extrapolation to 0 wt% Fe again yields the theoretical value for $L\beta/L\alpha$ of 0.5 (see above). At 2.5 kV, the small electron penetration depth of approximately 60 nm (Fig. 8) results in short path lengths for Fe L X-ray emission and practically no self-absorption. At 10 kV, the familiar quadrilateral appears, as seen in Figure 6, which widens at 15 and 20 kV (Fig. 10). However, there is no significant difference detectable between the 15 and 20 kV data, which means that the maximum sampling volume has been reached at 15 kV. Beyond 20 kV, no further widening of the quadrilateral is expected since the peak maxima do not shift with accelerating voltage anymore (Pavčević et al. 1989). It becomes clear from these findings that the flank method can be applied at 15 kV, which also provides the opportunity to simultaneously carry out conventional quantitative elemental analyses. This in turn will allow for synchronized chemical and Fe-valency mapping.

Self-absorption correction and quantification of Fe$^{3+}$ at 15 kV

In consequence, all garnets (Table 1) were remeasured under optimized conditions with the Jeol microprobe, i.e., at an accelerating voltage of 15 kV, 60 nA, a beam diameter defocused to 1 µm or focused to minimum diameter, and with the simultaneous measurement of $L\beta/L\alpha$ and major elements to take possible zonation into account. For the flank method, at least 10 repeated measurements were carried out on each sample with 300 s counting time both on the FeL$\beta$ and FeL$\alpha$ flank positions. The results are shown in Figure 11. The shape of the quadrilateral of Figure 6 is kept but it has increased in size as the result of increased self-absorption. This gives a better resolution in $L\beta/L\alpha$, which in turn will yield more precise ferric Fe values.

To quantify the relationship between Fe$^{2+}$ and $L\beta/L\alpha$, multiple linear regression was applied to a simple linear model, which takes both self-absorption effects into account, i.e., the correlation with Fe$^{2+}$ and the correlation with total Fe. In further tests, we checked for matrix effects from the chemical composition of the garnets. For multiple linear regression, Fe$^{2+}$ was taken as the dependent variable in a linear model:

- **Figure 8.** Sketch of the electron penetration depth in a garnet as a function of accelerating voltage based on a Monte Carlo simulation for almandine with a density of 4 g/cm$^3$. Fe$^L$ X-ray emission is increasingly subject to (self-) absorption with increasing accelerating voltage due to deeper X-ray generation depths and absorption path lengths through the sample.

- **Figure 9.** Dependency of the flank-method ratio on excitation energy. The difference between andradite and almandine increases between 2.5 and 15 kV. It does not change anymore at 20 kV compared to 15 kV. $E_c$ is the critical excitation energy for Fe$L\alpha$ X-ray emission.

- **Figure 10.** Variation of the flank-method ratios with total Fe for different excitation energies: 3 synthetic samples (And, Alm, Ski), and the McGuire almandine oxygen standard.
Fe\(^{2+}\) = a + b(Lβ/Lo)

(1)

Since the Fe\(^{2+}\)-Lβ/Lo dependency changes with variations in Fe\(^{3+}\) or total Fe content from the microprobe analysis, coefficients “a” and “b” in Equation 1 were formulated to be dependent on total Fe: a = a\(_1\) + a\(_2\)ΣFe and b = b\(_1\) + b\(_2\)ΣFe. Equation 1 can then be written as:

Fe\(^{2+}\) = A + B(Lβ/Lo) + CΣFe + DΣFe × (Lβ/Lo)

(2)

In a first step, the flank-method data from all the garnet analyses were used for the multiple regression, i.e., a total of 523 measurements. The coefficients A, B, C, and D are given in Table 2 together with the standard error of the estimate for Fe\(^{2+}\) in Figure 11 (the calculation of Fe\(^{3+}\)/ΣFe of approximately 4% in a sample with 10% total Fe).

In a second step, a reduced data set with 205 measurements was used. Measurements were omitted if: (1) the samples were highly inhomogeneous on the microscale (most almandine-spessartine and some grossular-andradite garnets); (2) the almandine-spessartine garnets contained small but varying amounts of Fe\(^{3+}\); (3) samples that were supposed to contain no Fe\(^{3+}\) but apparently did contain some. [They are from multiphase experiments and (b) those measurements were omitted where the samples were too inhomogeneous or the Fe\(^{2+}\) content was not correct (see text).

Notes

(a) 523 –9.61 (25) 18.59(26) –0.013(13) 0.160(10) 0.997 0.918
(b) 205 –10.11(17) 18.02(19) 0.015(8) 0.159(7) 1.0 0.386

TABLE 2. Results of multiple linear regression fits of the flank method measurements on synthetic garnets

<table>
<thead>
<tr>
<th>Fit</th>
<th>N</th>
<th>A</th>
<th>B</th>
<th>C</th>
<th>D</th>
<th>R</th>
<th>S</th>
</tr>
</thead>
<tbody>
<tr>
<td>(a)</td>
<td>523</td>
<td>–9.61(25)</td>
<td>18.59(26)</td>
<td>–0.013(13)</td>
<td>0.160(10)</td>
<td>0.997</td>
<td>0.918</td>
</tr>
<tr>
<td>(b)</td>
<td>205</td>
<td>–10.11(17)</td>
<td>18.02(19)</td>
<td>0.015(8)</td>
<td>0.159(7)</td>
<td>1.0</td>
<td>0.386</td>
</tr>
</tbody>
</table>

Notes: Fits (a) and (b) stem from the same data set at 15 kV: (a) contains all measurements; for (b), those measurements were omitted where the samples were too inhomogeneous or the Fe\(^{2+}\) content was not correct (see text). N = no. of measurements; A, B, C, and D = unstandardized coefficients according to Equation 2; R = correlation coefficient; S = standard error of the Fe\(^{2+}\) estimate in wt%. Parentheses show the error in the last digits (1sd).

Figure 12 shows the unstandardized residuals from the multiple linear regression fit for data set b. It demonstrates its very good quality and also shows that there are no matrix effects, at least for combinations of the elements Ca, Mg, and Al. The same conclusion also holds for Mn and Cr since the garnets with the highest contents of these elements show no resolvable deviation from the pyrope-grossular-almandine line (Fig. 11). The deviations seen for the Mn- and Cr-poor (high-Fe) garnets must be due to an Fe\(^{3+}\) component. The flank-method results are not affected by the chemical composition of these widely different garnets, i.e., there is no matrix effect. Consequently, Fe\(^{2+}\) can be calculated in wt% from a fit of data set b to Equation 2 that yields Fe\(^{2+}\) = –10.11(17) + 18.02(19) × (Lβ/Lo) + 0.015(8) × ΣFe + 0.159(7) × ΣFe × (Lβ/Lo) (3), with the numbers in brackets corresponding to the error in the last digits.

Solving Equation 3 for different Fe\(^{3+}\)/ΣFe ratios yields a set of Table 3.

\[
\frac{Lβ}{Lo} \approx \frac{a}{b} \left(1 + \frac{Lβ}{Lo} \right) \times \frac{1}{1 + \frac{a}{b}} \times (\frac{Lβ}{Lo})
\]

\[
Lβ = \frac{a}{b} \times Lo \times \frac{1}{1 + \frac{a}{b}} \times (\frac{Lβ}{Lo})
\]

FIGURE 11. Flank-method ratios obtained at 15 kV, 60 nA for all synthetic garnets plotted vs. the total Fe content. The overall ratio difference of Lβ/Lo between andradite and almandine has increased in comparison to the 10 kV data in Figure 6.
of curves for equal Fe speciation (Fig. 13). Fe$^{3+}$/ΣFe can now be directly estimated from the diagram, showing that most of the Mn- and Cr-containing garnets have Fe$^{3+}$/ΣFe up to 0.1. Also, samples that appeared inhomogeneous in terms of total Fe and the main elements are homogenous in Fe$^{3+}$/ΣFe since they plot along the lines in Figure 13.

**FIRST APPLICATION OF THE SELF-ABSORPTION CORRECTION**

**Natural garnets**

The flank method was established with synthetic garnets having a large variation in the total Fe and Fe$^{3+}$ contents. This enabled us to find systematic variations as well as to quantify Fe$^{3+}$/ΣFe by exploiting the effect of self-absorption. To test this method for natural samples, we extended our studies to homogeneous garnet megacrysts, five from the Biu and Jos Plateaus in Nigeria (Rankenburg et al. 2004), and one from Monastery ("MonB," courtesy A.B. Woodland). Their Fe$^{3+}$/ΣFe were previously determined by Mössbauer spectroscopy [A.B. Woodland in Rankenburg (2002) and pers. comm.]. Table 3 reproduces these data.

We obtained flank method data on these natural garnets and calculated Fe$^{3+}$/ΣFe with Equation 3. The results are compared with the Mössbauer data in Figure 14a. The Nigerian garnets with variable bulk Fe but very similar Fe$^{3+}$/ΣFe plot almost in one point along a 1:1 correlation line. This demonstrates the validity of the correction procedure for differing bulk Fe contents. The vertical 1σ error bars were derived from the mean of 10 flank-method measurements each. The Mössbauer error was taken as ± 0.01 for these garnets with a low Fe content. This value was suggested by Jayasuriya et al. (2004) for a good Mössbauer data set. Note that the error bars of the flank method and Mössbauer determination are similar in size. The almandine oxygen standard also plots on the 1:1 line at 0.04 Fe$^{3+}$/ΣFe, i.e., the Mössbauer value given by McGuire et al. (1992). The Monastery megacryst with significantly different Fe$^{3+}$/ΣFe plots off the 1:1 line. The two independent calibrations with the synthetic and with the natural garnets yield results for Mon B, which are not consistent with those for the

![Figure 12](image-url)

**Figure 12.** Residuals of the multiple linear regression for fit (b) of the synthetic garnets (Table 2) plotted vs. their chemical composition. There is no dependency on the chemical parameters.

![Figure 13](image-url)

**Figure 13.** Equilines of Fe$^{3+}$/ΣFe calculated from Equation 3 in a diagram of $Lβ/Lα$ vs. the total Fe content in the synthetic garnets. The full symbols represent the garnet compositions used to derive Equation 3. From a comparison with the equilines, the Fe$^{3+}$/ΣFe content of the garnets with open symbols may be estimated.
The megacrysts from Nigeria are of gem quality. The megacrysts from Nigeria are of gem quality. The megacrysts from Nigeria are of gem quality. The megacrysts from Nigeria are of gem quality. Proceeding further, we used two Nigerian megacrysts (Ker GB and Damknolle) and the McGuire almandine with widely differing total Fe to obtain a calibration with fit parameters derived solely from these three natural garnets. We then calculated Fe\(^{3+}\)/ΣFe for the other megacrysts and compared the results with Mössbauer data in Figure 14b. There is practically no difference between Figures 14a and 14b except for somewhat larger error bars. There is practically no difference between Figures 14a and 14b except for somewhat larger error bars in Figure 14b. The flank method again yields a significantly higher Fe\(^{3+}\)/ΣFe than the Mössbauer value for Mon B.

Modification: Linear self-absorption correction for natural garnets

Figure 13 shows gently curved equilines of Fe\(^{3+}\)/ΣFe that reproduce the compositions for the synthetic and the natural garnets quite successfully. The composition of natural garnets mostly lies between 5 and 20 wt% total Fe. A closer look at this range shows that the data may be equally well approximated by a straight line (compare the solid and dashed lines in Fig. 15a). In this figure, only the Fe\(^{2+}\)-garnets shown with full symbols were used for the linear regression calculation (with only one fit parameter) to derive the line for Fe\(^{3+}\)/ΣFe = 0. This makes Fe\(^{2+}\) dependent only on the \(L_b/L_o\) ratio as in Equation 1. For the calibration shown in Figure 15b, the same three natural garnets as in Figure 14b, along with their Fe\(^{3+}\) content from Mössbauer analysis, were used to derive the Fe\(^{3+}\)/ΣFe = 0 line. The difference between this line and the fit with the synthetic garnets is barely distinguishable, indicating that the calibrations made with synthetic garnets and natural garnets yield basically the same results.

This result leads to a simplified self-absorption correction suitable for the restricted compositional range of natural garnets. Calibration is performed using the above 3 natural garnets. Plotting the measured \(L_b/L_o\) ratio vs. the Fe\(^{2+}\) content obtained from Mössbauer spectroscopy also gives a linear relationship (Fig. 16a). The \(L_b/L_o\) ratio of a garnet with unknown Fe\(^{2+}\) but known total Fe content will plot below the regression line if the abscissa is taken as bulk Fe (insert in Fig. 16a). The magnitude of the deviation from the calibration curve (Δ ratio) is a measure of the Fe\(^{2+}\) content of the sample. It is calculated as the difference between the corresponding point on the calibration curve and the measured \(L_b/L_o\) ratio. Plotting the Δ ratios of the garnets used for calibration against their Mössbauer Fe\(^{2+}\) contents also yields a linear relationship (full symbols in Fig. 16b). The open symbols are the garnets taken as unknowns. Their Fe\(^{2+}\) (wt%) can be read off graphically or calculated with the calibration line.

The results obtained in this way are compared in Figure 17 with their respective Fe\(^{3+}\)/ΣFe Mössbauer values. The error bars correspond to the 1σ uncertainty for the mean value of 10 to 20 analyzed spots in the sample (see text and Table 3). The close similarity demonstrates the validity of the approach and that a restricted number of well-characterized natural garnet standards may be used in the future.


data from Table 3:

<table>
<thead>
<tr>
<th>Sample</th>
<th>(\Sigma Fe) (wt%)</th>
<th>Fe(^{3+})/ΣFe Mössbauer</th>
<th>Fe(^{3+})/ΣFe Flank Method</th>
<th>1σ Flank Method</th>
<th>N</th>
<th>Reference for sample details</th>
</tr>
</thead>
<tbody>
<tr>
<td>Alm McGuire</td>
<td>20.21(6)</td>
<td>0.04</td>
<td>calibration</td>
<td>–</td>
<td>16</td>
<td>McGuire et al. 1992</td>
</tr>
<tr>
<td>Damknolle</td>
<td>15.18(3)</td>
<td>0.051</td>
<td>calibration</td>
<td>–</td>
<td>10</td>
<td>Rankenburg et al. 2004</td>
</tr>
<tr>
<td>Ker GB</td>
<td>7.54(2)</td>
<td>0.048</td>
<td>calibration</td>
<td>–</td>
<td>10</td>
<td>Rankenburg et al. 2004</td>
</tr>
<tr>
<td>Mir 13</td>
<td>10.66(3)</td>
<td>0.050</td>
<td>0.049</td>
<td>0.010</td>
<td>20</td>
<td>Rankenburg et al. 2004</td>
</tr>
<tr>
<td>Mir 2</td>
<td>10.04(2)</td>
<td>0.051</td>
<td>0.025</td>
<td>0.021</td>
<td>10</td>
<td>Rankenburg et al. 2004</td>
</tr>
<tr>
<td>Mir 23</td>
<td>8.31(2)</td>
<td>0.053</td>
<td>0.089</td>
<td>0.018</td>
<td>10</td>
<td>Rankenburg et al. 2004</td>
</tr>
<tr>
<td>Mon B</td>
<td>8.89(4)</td>
<td>0.115</td>
<td>0.192</td>
<td>0.020</td>
<td>10</td>
<td>Woodland, pers. comm.</td>
</tr>
</tbody>
</table>

Note: The error in Fe\(^{3+}\)/ΣFe (Mössbauer) is approximately ±0.01; the Mössbauer data of Rankenburg et al. are from A.B. Woodland; N = number of measurements; brackets show the error in the last digit.
by the use of higher beam currents, longer counting times, and/or an increased number of repeated measurements on homogeneous samples. A crucial point is the scarcity of good Fe\(^{2+}/Fe^{3+}\) standards with a very restricted range in Fe\(^{3+}/Fe^{2+}\) ratio that are unambiguously related to the valence state of Fe via the differences in self-absorption of Fe\(^{2+}\) and Fe\(^{3+}\). Our findings are: (1) The flank method Fe\(L_\beta/L_\alpha\) ratios systematically vary with Fe\(^{3+}/Fe^{2+}\); (2) These variations are related to differences in the self-absorption of Fe\(^{2+}\) and Fe\(^{3+}\); (3) The total self-absorption is a linear superposition of both effects and may be quantified with multiple linear regression; (4) The chemical composition does not measurably affect the \(L_\beta/L_\alpha\) ratio of the studied garnets i.e., there are no matrix effects.

Matrix correction is a key feature of conventional quantitative EPMA analysis. For the flank method, the data are not corrected for matrix effects because the measurement position for Fe\(L_\alpha\) exactly lies on the \(L_\alpha\) absorption edge (Fig. 1) where the size of matrix absorption coefficients are most uncertain. Empirically we have shown that material effects may be neglected (see above) and this is also feasible on theoretical grounds. It is known that photoelectric absorption is very strong if the photon energy is close to the absorption energy. In the case of Fe\(L_\alpha\)-lines with 2p ↔ 3d electron transitions, the 2p3d X-ray emission energy is close to the 2p3d X-ray excitation (absorption) energy (Fig. 18). In such a situation X-ray absorption and X-ray emission processes do occur coherently leading to resonant X-ray processes (de Groot 2001). As a result, self-absorption will be very strong. For Fe\(L_\alpha\) X-ray emission, electron transitions take place from a filled 3d electron state to a 2p core hole. The X-ray absorption process is roughly the opposite (Fig. 18), i.e., it excites a 2p core electron into the empty 3d shell. In a ligand-field, the 3d levels split according to the site symmetry of the Fe atom, as shown schematically. Depending on the ground-state configuration of Fe\(^{2+}\) and Fe\(^{3+}\), this leads to a different spectral shape for both the X-ray absorption and the emission spectra. In the garnet structure, Fe\(^{2+}\) and Fe\(^{3+}\) occupy the dodecahedral and octahedral sites, respectively. Therefore, the crystal-field splitting is mainly influenced by Fe\(^{2+}/Fe^{3+}\). Consequently, the shape and fine structure of the \(L_\alpha\) absorption edge changes significantly with the Fe oxidation state (see Fig. 1; see also van Aken and Liebscher 2002). This leads to a shift of the energy of the absorption maximum and thus to selective self-absorption of the Fe species. As argued above,
self-absorption is enhanced due to resonant X-ray processes. The resonance is presumably so strong that matrix absorption by chemical components other than Fe becomes negligible. This explains the independence of the flank method ratios from chemical composition.

The flank method enables the quantitative determination of $\text{Fe}^{3+}/\Sigma\text{Fe}$ with the electron microprobe within typical microprobe sampling volumes. It could be shown for garnets that the dominant component influencing the $\text{Fe}L\alpha$ and $\text{Fe}L\beta$ intensity ratios is the selective self-absorption of the FeL X-ray emission by $\text{Fe}^{2+}$ and $\text{Fe}^{3+}$. Major element composition does not play a role so that raw intensities uncorrected for matrix absorption may be used for calculations. The $L\beta/L\alpha$ ratios determined with the flank method appear sensitive to the coordination polyhedron for Fe, i.e., to the electronic surroundings of the Fe ions.

This may be seen from previous work of Höfer et al. (1994, 2000) who presented 10 kV flank measurements on natural olivines, synthetic wüstites, and synthetic Fe-rich spinels, and from Enders et al. (2000) who applied the same technique to natural amphiboles. These data have been replotted in Figure 19 along with our 10 kV garnet data and previously unpublished data from our work on pyroxenes, hematite, and ilmenite to demonstrate the effect of self-absorption and coordination polyhedra. For all mineral groups, except spinel, linear arrays originate with differing slopes that intersect the y-axis at around 0.5 for $L\beta/L\alpha$. The slopes of these lines appear to be related to the coordination polyhedra of $\text{Fe}^{2+}$. That is, garnets with dodecahedral $\text{Fe}^{2+}$ show the highest slope. Amphiboles with slightly distorted eight- and sixfold coordinated polyhedra for $\text{Fe}^{3+}$ have a lower slope. Pyroxenes with more tight eight- and sixfold polyhedra and olivines with only sixfold polyhedra have a still lower slope. Wüstites with only sixfold coordination have a different y-intercept but their slope corresponds to the pyroxene, olivine, and hematite structure minerals. Spinels with octahedral to tetrahedral coordination plot at the lowest $L\beta/L\alpha$ values.

Measurement of the flank-method ratios can be very precise so that the limiting factor for accurate determination of $\text{Fe}^{3+}/\Sigma\text{Fe}$ is the quality of the standards. They need to be homogeneous and well defined. The potential of the flank method has been shown in this study on mantle garnets where we found an error of only $\pm 0.02$ for $\text{Fe}^{3+}/\Sigma\text{Fe}$ in samples with 8–11 wt% total Fe.

The search for suitable and accurate $\text{Fe}^{3+}/\Sigma\text{Fe}$ standards will be one main goal in future work to improve the accuracy of $\text{Fe}^{3+}$ determination by electron microprobe. Another goal will be to establish the flank method in minerals other than garnet. We are currently extending our investigations to synthetic pyroxenes (Höfer et al. 2004). The main advantage of the flank method over Mössbauer spectroscopy, and also other methods such as EELS or XANES, is that measurements are possible within 10 min on a micrometer scale with the same or better accuracy in samples with low Fe contents. A further advantage is that sample preparation is easy and that electron microprobes are readily available in many institutions throughout the world.
ACKNOWLEDGMENTS

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HÖFER AND BREY: IRON OXIDATION STATE BY ELECTRON MICROPROBE

884

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APPENDIX

Summary of measurement procedure

The following gives a practical overview on how to use the flank method with either a Cameca or a Jeol microprobe.

(a) Spectrometer calibration (daily procedure)

1. 25 kV, 80 nA: peak search for FeKα (9th order) on Fe metal (Integral mode; optimize SCA (PHA) conditions of the TAP spectrometer to be used for the flank method measurements for this line). For best resolution, always use the smallest available detector slit for this spectrometer (300 μm for Jeol JXA 8900).

2. Quick qualitative scan over this peak for better definition of peak maximum (2–3 min).

3. Define the FeKα (9th order) on Fe metal for the Cameca microprobe after (following Höfer et al. 2000, cf. p. 66–67 and Fig. 3). For our Jeol microprobe, we have written a program “PeakFitH,” which automatically runs the routine (4) below. This program can be obtained free of charge from the first author.

4. Run the program PeakFitH for the TAP spectrometer with the newly defined value for the FeKα (9th order) peak maximum: 3 repetitions at 6 spectrometer positions are sufficient, counting time per position: 100 s (Jeol; 10 min); or accumulation to a certain number of counts (see Höfer et al. 2000). The 6 measuring positions (“window” values, i.e., the symmetrical distances of fe1 and fe2 around the peak maximum, see Fig. 3 in Höfer et al. 2000) may have to be adjusted to your spectrometer (our values for Jeol: window 52 μm, shift 14 μm; for Cameca: window 25 sinΘ steps, shift 5 sinΘ steps).

5. Transfer the results (intensities and beam current values) into an Excel spreadsheet; calculate beam current-normalized ratios fe1/fe2 (according to Höfer et al. 2000).

6. Plot the 3 values for fe1/fe2 vs. the spectrometer position, calculate the corresponding regression line and the spectrometer shift for fe1/fe2 = 1 for the current measuring day. Note: fe1/fe2 = 1 if the peak is symmetrical in this region, otherwise take another value suitable for your spectrometer. This value also represents an external energy calibration if it is set equal to the theoretical value of FeKα (9th order). The fe1/fe2 value is specific for every spectrometer and microprobe, and should be maintained as an internal reference and as energy calibration for all flank-method measurements.

7. This complete procedure takes about 30 min.

(b) Flank-method measurement

1. After the spectrometer calibration, set the microprobe to 15 kV, 60 nA.

2. Create the measurement condition file. For a combined flank method and quantitative analysis, load a suitable quantitative measurement program. You may increase counting times for the elements to obtain the best quantitative results since one TAP spectrometer will be busy carrying out flank measurements for 10 minutes. Add two “fake” elements for the flank method measurements, e.g., As for FeLα and Br for FeLβ. Insert the TAP fake “peak” position for the flank method measurements by adjusting it with the spectrometer shift value from the above calibration procedure (6) of the day. The fake “peak” measurement positions for FeLα and FeLβ have to be determined only once from difference spectra (Fig. 1) when you begin flank measurements with your microprobe. Please ensure that you have gone through the spectrometer calibration procedure beforehand. Allow long measuring times for FeLα and FeLβ such as 300 s each, no background measurements, differential mode, narrow energy window to avoid overlaps with higher-order X-ray emission lines, and optimized SCA (PHA) conditions for FeLα on Fe metal.

3. The statistically more precise method is to set a certain number of counts to be accumulated (Höfer et al. 1994). However, our Jeol microprobe does not allow this for low count rates at present. We therefore resorted to 300 s counting times on each peak. To avoid contamination artifacts, measure Lβ first and then Lα on the same spectrometer. We also tested longer measuring times of 600 s each that are free of contamination effects. To increase counting statistics, we suggest increasing the number of measurements but not necessarily the counting time. If your sample is homogeneous and you wish to increase the total number of points, you may choose measurement times shorter than 300 s on single points and average the data. But the longer 10 min measurement times (300 s on each position) have the advantage that good quantitative results are achieved simultaneously with the Fe³⁺/ΣFe determination.

4. Input the measured points by using the usual quantitative measurement program. Begin by measuring an almandine-rich sample as a reference to monitor the quality of the flank-method measurement. One measurement for this is enough. Then measure 2–3 standard samples with different Fe³⁺/ΣFe and/or total Fe contents. Next, measure your samples (preferably also 10 times or more on different spots). Insert a few measurements on almandine every now and then to check the stability of the spectrometer.

5. Run the program routinely.

6. For quantitative determination, run an off-line matrix correction by kicking out the 2 “fake” elements, As and Br.

7. For flank-method ratios, run a summary with net intensities of only As and Br. Add the beam current values to the summary. Insert these data into an Excel spreadsheet. Multiply the net intensities (cps/μA) by the beam current to obtain count rates and calculate the flank method ratios.

8. These data can be used for Fe³⁺/ΣFe determination.